

Thermodynamic Evaluation of a Refrigeration Bottoming Cycle on the Efficiency of a Condensing Supercritical Carbon Dioxide Power Cycle

Phalgun N Malupillai
Graduate Researcher
University of Washington
Seattle, Washington

Jon D. McWhirter, PhD, P.E.
Sr. Manager MCFR, Thermal Hydraulics
TerraPower LLC,
Bellevue, Washington



Phalgun Malupillai is a Mechanical Engineering MS student at the University of Washington, currently working as a research intern with TerraPower. He holds a B.Tech in Chemical Engineering from NIT Trichy, and has a strong interest towards thermodynamics, clean energy, and carbon capture.



Jon McWhirter is Senior Manager, Molten Chloride Fast Reactor Thermal Hydraulics at TerraPower LLC in Bellevue, WA. He holds a Doctorate from The University of Texas at Austin, and Professional Engineer licenses in the states of Idaho and Washington.

Nomenclature

ΔT	Temperature Difference ($^{\circ}C$)	Q	Heat Duty (kJ)
COP	Coefficient of Performance	s	Specific Entropy (kJ/kg-K)
C_p	Specific Heat Capacity (kJ/kg-K)	sCO_2	Supercritical CO ₂
HEX	Heat Exchanger	T	Temperature ($^{\circ}C$)
HTR	High Temperature Recuperator	T_o	Dead state Temperature ($^{\circ}C$)
i	Irreversibility (kJ) or (kJ/s)	T_{Rej}	Heat Rejection Temperature ($^{\circ}C$)
LTR	Low Temperature Recuperator	T-S	Temperature - Entropy
\dot{m}	Mass flow rate (kg/s)	η	Efficiency
MSR	Molten Salt Reactor	Φ	Availability / Exergy (kJ)

9 Abstract

10 Supercritical Carbon Dioxide (sCO₂) power cycles have recently garnered attention for use with high
11 temperature Gen IV nuclear reactors such as Molten Salt Reactors (MSR) due to prospects for improved
12 performance in high temperature regions, with the Recompression model being one of the most studied
13 configurations where the heat rejection step occurs near the CO₂ critical point of 30.98°C. An alternative
14 well researched variant is the transcritical condensing CO₂ cycle which reduces irreversibility in heat
15 rejection and shows higher efficiencies and specific work produced. The stable operation of an MSR
16 benefits from maintaining steady operating points, heat transfer and approach temperature between the
17 salt loop and the power cycle. Given the proximity of the critical point of CO₂ to ambient temperatures,
18 and the variability of thermodynamic properties near this point, sCO₂ cycles are susceptible to
19 unwelcome transients associated with ambient temperature fluctuations. Further, condensing cycles,
20 despite showing promising performance, are limited to regions with consistent ambient temperatures
21 which allow heat rejection below the critical temperature, year-round.

22 This study proposes a “Thermal Adapter” model, which uses a refrigeration bottoming cycle to maintain
23 stable operating points in a condensing sCO₂ recompression cycle, while allowing condensation
24 irrespective of ambient conditions, thereby addressing both these concerns. The isothermal heat
25 rejection from the CO₂ (in contrast to a Brayton cycle) to refrigerant, and refrigerant to the environment
26 mitigates some of the work lost through the operation of the refrigeration cycle. A thermodynamic and
27 exergy analysis was conducted comparing the Thermal Adapter model to the standard recompression
28 cycle at varying environmental temperatures. The results indicate that the Thermal Adapter successfully
29 isolates the CO₂ cycle and nuclear operation from transient ambient conditions, exhibiting a greater
30 specific work at all ambient temperatures, a higher net efficiency at $T_{Rej} > 32^{\circ}\text{C}$, and operational simplicity
31 on comparison with the Reference model.

32 Introduction

33 Recent years have seen large investments and technological advancements made towards the launch of
34 Gen IV nuclear reactors, as a step towards more efficient, and inherently safer nuclear power. These
35 reactors generate heat at higher temperatures compared to most existing powerplants, with
36 TerraPower’s Molten Salt Reactor slated to operate near 650°C, allowing for a higher Carnot efficiency in
37 the coupled power cycle.

38 One such advanced power cycle poised towards high temperature application is the supercritical carbon
39 dioxide (sCO₂) cycles which operates at high efficiencies with relatively compact and simple plant layouts,
40 providing economic and operational benefits over competing supercritical steam cycles [1], and helium
41 Brayton cycles [2]. The advantages of sCO₂ cycles stem from the low compressibility of CO₂ near the
42 critical point, minimizing compression work, ability to internally recuperate heat, and the high density of
43 CO₂ in the expansion process, allowing for smaller machinery, and higher energy density [3]. The low
44 critical point of CO₂ at 31°C and 7.38 MPa allow the cycle to reject heat near ambient conditions. The
45 recompression cycle which operates entirely above the critical pressure is one of the most studied
46 configurations [4]. The condensing transcritical recompression CO₂ cycle has been well studied by [5]
47 and involves condensation of a part of the CO₂ stream, which is pumped from the liquid phase to
48 supercritical pressure, resulting in a higher net power and efficiency [6]. Wright et al. have described
49 how a slight reduction in heat rejection temperature can lead to significant efficiency gains [7]. The

50 higher efficiency of the condensing cycle can be partially explained by the slightly lower heat rejection
51 temperature implemented to facilitate condensation, resulting in a higher Carnot efficiency. The gain in
52 the real cycle efficiency is greater than the increase in the Carnot efficiency associated with the wider
53 temperature range. This observation is explained through the greater second law efficiency of the cycle
54 arising from nearly isothermal heat rejection via condensation and thereby lower irreversibility.

55 The common concerns associated with the industrial implementation of sCO₂ cycles are the variations of
56 the thermodynamic fluid properties in the vicinity of the critical point and the impact this may have on
57 the overall power cycle. Unlike traditional steam cycles, CO₂ cycles reject heat near ambient conditions,
58 and given the critical point of CO₂ is near the ambient temperature in most geographies, environmental
59 temperature fluctuations may further disrupt the stable operation of a sCO₂ power cycle. The steady
60 performance and operation of molten salt reactors is benefitted by a constant, and predictable heat
61 transfer out of the salt loop and into the CO₂ cycle, hence such aforementioned disturbances in the CO₂
62 cycle as a result of transient ambient conditions are undesirable. The condensing sCO₂ cycle solves one
63 facet of this problem, by rejecting heat from the cycle at temperatures significantly below the critical
64 temperature. This process bypasses the critical point, and thereby any associated thermodynamic
65 instability regions, while exhibiting appreciable gains in efficiency and specific work. Nonetheless, such
66 cycles are limited geographically to regions where a sufficiently cold natural heat sink is available and are
67 yet susceptible to variations in external temperatures.

68 This study proposes the use of a refrigerating bottoming loop as a “Thermal Adapter”, which couples
69 with the heat rejection step of a condensing sCO₂ cycle. The Thermal Adapter will serve to maintain a
70 constant interface temperature with the CO₂, thereby isolating the power cycle and nuclear operations
71 from variations in external temperatures, while allowing for condensation of CO₂ irrespective to climatic
72 conditions. A well-designed refrigeration loop will employ heat transfer between both the CO₂ cycle, and
73 the environment through largely isothermal phase change processes, thereby minimizing irreversibilities,
74 and reaping the benefits of a condensing cycle, which recoups some of the work lost through the
75 introduction of an additional process. This study reports the thermodynamic and exergy analyses on the
76 implementation of a Thermal Adapter and outlines the use cases wherein the use of such a configuration
77 results in net efficiency and performance benefits.

78

79 **Modelling**

80 The thermodynamic modelling of the sCO₂ power cycle is carried out through the Aspen HYSYS modelling
81 software, using the Peng-Robinson equation of state. The alternate equations of state considered were
82 Lee-Kesler-Plocker and Redlich-Kwong-Soave. Peng-Robinson was selected due to the ease of
83 computation, and accuracy near the critical point.

84 The accuracy of the Aspen HYSYS model was verified through the replication of the state points in the
85 recompression cycle described by Wright et al. [7]. The model showed a variance of less than 1% in the
86 prediction of efficiency and a variance of 3.6% in prediction of specific power.

87 **Reference Model**

88 The Reference model is based on the well-studied recompression Brayton cycle model and is modelled
89 around the schematic described by Wright et al. [7], with some modifications as documented in Table 1.

90 The heat source is defined as an isothermal reservoir at 650°C. This analysis studies the effect of ambient
91 temperature fluctuation on the sCO₂ power cycle, hence the heat rejection temperature is varied from
92 15°C – 45°C. The high pressure point is 20 MPa, and the low pressure point post expansion is set slightly
93 above the critical pressure at 7.7 MPa, when the rejection temperature is above the critical temperature
94 of CO₂.

95 For a fair comparison with the Thermal Adapter model, the reference model is designed to allow for
96 condensation when the rejection temperature is sufficiently below the critical temperature ($T_{Rej} < 30^{\circ}\text{C}$).
97 This study does not directly refer to the ambient temperature as local conditions such as availability of
98 running water, the humidity of air, etc. will affect the approach temperature between the power cycle fluid
99 and ambient conditions, therefore the rejection temperature is used to denote the temperature of the
100 working fluid post the heat rejection step. In the Reference model, T_{Rej} refers to CO₂ temperature exiting
101 the rejection step. The turbine outlet pressure is dropped to the saturation pressure corresponding to the
102 heat rejection temperature, allowing for condensation during the heat rejection step. Studies have shown
103 that the low variation of density between the vapor and liquid phases of CO₂ allow for a compressor to
104 function as a pump for liquid CO₂, albeit at lower efficiencies. This study credits the Reference model with
105 the ability to pump condensed CO₂ by means of a compressor, with no penalty on adiabatic efficiency.

106 **Thermal Adapter Model**

107 The previously defined reference model has been modified to include a simple propane refrigeration cycle
108 which facilitates the condensation of the CO₂ in the CO₂ heat rejection step. The closed propane loop
109 consists of a compressor, compressing the propane to the saturation pressure at ambient conditions. The
110 propane vapor rejects heat to the environment, condensing into a saturated liquid, which is throttled to
111 form a two-phase mixture. The throttle outlet pressure is based on the level of Joule-Thompson throttling
112 required to drop the temperature to 12°C. The propane evaporates in the CO₂-Propane phase change heat
113 exchanger consequently condensing CO₂ at 15°C. The low temperature range of the propane cycle (12°C
114 to T_{Rej}), results in a high COP cycle. T_{Rej} in this model refers to the temperature of propane leaving the
115 condenser. The propane cycle effectively isolates the CO₂ loop from any fluctuations in the ambient
116 temperature through modifications in the mass flow rate, and pressure ratio to maintain a constant CO₂-
117 propane interface temperature. Figure 1 displays a process flow diagram highlighting the thermodynamics

118 state points for a $T_{Rej} = 40^\circ\text{C}$ case. The theoretical total efficiency of such a combined cycle can be calculated
119 as follows.

$$120 \quad \eta_{Overall} = \eta_{Top} - \frac{(1 - \eta_{Top})}{COP}$$

121 The Thermal Adapter leverages the CO₂ condensation to replace the main compressor with a centrifugal
122 pump. The condensation of CO₂ allows the turbine to operate at a higher pressure ratio, expanding the
123 fluid to the expected saturation pressure at 15°C, well below the critical pressure to which traditional
124 cycles are limited. Further, the constant heat rejection temperature of 15°C allows for finetuning of the
125 split flow ratio to allow for more effective heat transfer in the low temperature recuperator.

126 Split Flow Ratio

127 The split flow ratio is defined as the amount of flow directed to the heat rejection step, and consequently
128 the pump and LTR, as opposed to the re-compressor. This ratio is dictated by the heat transfer in the LTR.
129 The hot side carries the low-pressure vapor CO₂, while the cold side carries the high-pressure liquid CO₂
130 which undergoes evaporation and superheating in the LTR. The difference in thermodynamic conditions
131 results in a high difference in the specific heat capacity between the streams potentially leading to
132 unfavorable heat transfer [8]. To allow for a constant approach temperature across the heat exchanger,
133 the total heat capacity of both flows must be matched.

$$134 \quad \dot{m}_h * Cp_h * \Delta T_h = \dot{m}_c * Cp_c * \Delta T_c$$

$$135 \quad \text{For constant approach: } \Delta T_h = \Delta T_c ; (\dot{m}_h * Cp_h - \dot{m}_c * Cp_c) \rightarrow 0$$

136 Given the specific heat capacity is fixed by the pressure ratios, the mass flow is then controlled through
137 the split flow ratio to optimize operation. Figure 2 displays the effect of split flow ratio on the difference
138 in mass capacities ($\dot{m}_h C_{p,h} - \dot{m}_c C_{p,c}$) across the heat exchanger and the corresponding impact on overall
139 cycle efficiency. The peak cycle efficiency is found at a split flow ratio of 0.544.
140

141 **Results and Discussion**

142 The performance of the Thermal Adapter Model and the Reference model are evaluated and compared
143 under varying ambient temperature conditions (heat rejection temperature varies from 15°C to 45°C). The
144 effect of the propane refrigeration loop on overall system performance, and the ability to isolate the power
145 cycle from extraneous thermal fluctuations is documented in this section. The primary parameters of
146 concern are system efficiency, specific work, and CO₂ inlet temperature to the primary salt heater.

147 Efficiency

148 Figure 3 shows that the thermal adapter model exhibits a significantly higher efficiency at all rejection
149 temperatures above the critical point ($T_{Rej} > 33^\circ\text{C}$). At points greater than the critical temperature, the
150 Thermal Adapter model leverages isothermal phase change process to transfer heat from both the CO₂
151 into the propane loop and reject heat from the propane loop to the environment. This results in minimal
152 heat transfer related entropy gain, and consequently high efficiencies.

153 At rejection temperatures below the critical temperatures, the theoretical reference model reaps the
154 benefits of a condensing transcritical CO₂ cycle, without the parasitic load of the propane refrigeration
155 cycle, resulting in a theoretical higher efficiency. However, to reject heat via condensation, the system
156 must operate at higher pressure ratios, and repurpose the compressor to pump liquid CO₂. Studies have
157 shown that such an application of compressors is possible, albeit at a less than ideal efficiency point
158 thereby diminishing the overall cycle efficiency [9]. For the reference model to outperform the thermal
159 adapter model while condensing CO₂ ($T_{Rej} < 31^{\circ}\text{C}$), a compressor designed for an 85% peak efficiency for
160 vapor phase application, must exhibit an adiabatic efficiency of greater than 75% while pressurizing the
161 liquid phase. A more feasible operational strategy under colder ambient conditions ($T_{ambient} < 31^{\circ}\text{C}$) in a
162 standard sCO₂ cycle would be to restrict the cooling of low-pressure CO₂ to near the critical temperature,
163 consequently underleveraging the available natural heatsink, given the restrictions imposed by the
164 turbomachinery. The operation of the Thermal Adapter model is not restricted by the ambient
165 temperatures to the same degree, however, is still susceptible to efficiency losses arising from off-design
166 operation of the Propane compressor.

167 Specific Work

168 Figure 4 depicts the variation of specific work produced by the cycle, with variations in the heat rejection
169 temperature. The improved performance of the thermal adapter model is attributed to the higher-
170 pressure ratio across the turbine, and the lower work required to pump liquid CO₂. The adapter model
171 exhibits a 15% higher net work at T_{Rej} of 26°C, quickly rising to a 60% higher net work at a T_{Rej} of 39°C when
172 compared to the reference model.

173 Impact on Salt Heat Exchange Process

174 The primary heat input step in the proposed nuclear power application is through a salt heat exchanger.
175 The stable operation of the nuclear process benefits from maintaining a steady heat transfer between the
176 salt and the attached power cycle. Figure 5 and Figure 6 depict the impact of variations in environmental
177 conditions on the entry temperatures of CO₂ into the salt heat exchanger, as well as the heat absorbed by
178 the power cycle for a 1 kg/s mass flow rate of CO₂, respectively. The graphs clearly show that the propane
179 refrigeration loop successfully isolates the power cycle, and consequently the nuclear process from any
180 transient environmental conditions, providing a constant CO₂ inlet temperature to, and constant Q_{out} from
181 the salt loop.

182 T-S Diagram at $T_{Rej} = 40^{\circ}\text{C}$

183 A comparison between the T-S diagrams of the Thermal Adapter model, and the Reference model, when
184 operating under identical temperature boundary conditions is depicted in Figure 7 and Figure 8. This chart
185 depicts the ability of the propane refrigeration loop to pull the CO₂ into the liquid phase, allowing for near
186 isothermal heat rejection, as opposed to the large temperature gradient across the rejection step in the
187 reference model. The ability of the propane loop to exchange heat isothermally with both the CO₂ loop,
188 and the environment, results in an overall more efficient system. This behavior is consistent in all cases
189 where $T_{Rej} > 32^{\circ}\text{C}$.

190 **Exergy Analysis and Second Law Efficiency**

191 A detailed exergy analysis was conducted to thermodynamically analyze and compare the performance of
192 the Thermal Adapter model with the Reference model. The analysis assumes heat input through an
193 isothermal heat source at 650°C, and an isothermal heat sink at the rejection temperature. The exergy
194 entering the system, and the irreversibility of each of the unit processes are calculated based on the
195 formulae in Table 2.

196 Figure 9 and Figure 10 display the comparison of irreversibility as a percentage of exergy entering the
197 system for a 40°C heat rejection case. The irreversibility of the “Rejection” process in the Reference model
198 refers solely to the work lost while rejecting heat to the environment. The irreversibility of rejection in the
199 Thermal Adapter model refers to the sum of the work lost across the CO₂-propane heat exchanger, along
200 with the work lost in the propane cycle (compressor, condenser, and throttle).

201 The data shows that the irreversibility of heat rejection is significantly lower in the Thermal Adapter model,
202 despite the presence of additional equipment and processes. The facilitation of isothermal heat rejection
203 from the CO₂ loop to the Propane loop as opposed to direct heat rejection to the atmosphere results in an
204 overall less irreversible (more efficient) system. This behavior is exhibited for all cases where the heat
205 rejection temperature is greater than the critical temperature of CO₂.

206 **Conclusion**

207 The benefits of a condensing transcritical CO₂ power cycles are well established in literature. A primary
208 barrier associated with such a process is the proximity of the critical temperature to ambient conditions,
209 wherein variations in environmental temperatures disrupt the ability to condense CO₂ in the heat rejection
210 step. Further, this study is directed towards the application of sCO₂ cycles in the nuclear industry, where
211 maintaining stable operating conditions in the power cycle and the nuclear plant, irrespective of transient
212 ambient conditions, are of importance. This analysis follows a thorough modelling effort resulting in the
213 development of a “Thermal Adapter” system which incorporates a propane refrigeration bottoming loop
214 for a condensing CO₂ cycle and evaluates the ability of a such a loop to isolate the power cycle from
215 ambient temperature fluctuations, and the effect of the bottoming loop on the thermodynamic
216 performance of the overall system. The proposed system is compared against a standard condensing sCO₂
217 power cycle rejecting heat directly to the environment.

218 A thermodynamic and exergy analysis resulted in the following conclusions:

219 • The Thermal adapter model successfully isolates the power cycle from ambient thermal
220 fluctuations, as measured by the constant inlet temperature to, and heat duty absorbed from the
221 salt HEX. In contrast, a 25°C variation in ambient temperature in the reference model results in an
222 83°C variation in CO₂ inlet temp to the salt HEX.

223

224 • The Thermal Adapter model boasts of a higher specific work at almost all ambient temperatures,
225 despite the presence of additional rotating equipment (12% higher at a T_{Rej} of 25°C, and 61% higher
226 at 40°C). This is attributed to a higher pressure ratio across the turbine, and ease of pressurizing
227 liquid CO₂ as opposed to a gas compression process.

228

- 229 • The Thermal Adapter model exhibits a significantly higher efficiency than the reference model at
230 all $T_{Rej} > 32^{\circ}\text{C}$, due to the isothermal heat rejection via condensation. The exergy analysis around
231 this process shows that this isothermal rejection step compensates for the irreversibilities added
232 by the propane loop, resulting in a net efficiency gain when compared to the Reference model.

233 The reference model in this study shows a higher efficiency when the T_{Rej} is between 25°C and 32°C , as the
234 model benefits from the condensation of the CO₂ without the parasitic load of the refrigeration cycle.
235 However, this performance is contingent on the compressor pressurizing liquid CO₂ at the same efficiency
236 as gaseous CO₂, and such high performance has been shown to be unlikely. Further, the thermal adapter
237 model may be modified such that when the ambient conditions are sufficiently low, the propane loop may
238 be shut down, and the CO₂ stream can be redirected to reject heat directly to the environment, allowing
239 for natural condensation. Following these two considerations, the authors believe that the Thermal
240 Adapter model may thermodynamically outperform the reference model at all environmental conditions;
241 however, further investigation is recommended. Additional perks of the newly devised set up are the
242 ability to design a singular power cycle for application in various climatic zones with minor modifications
243 to the refrigeration loop alone, which result in engineering design cost saving and commonality of
244 hardware.

245

246 **References**

- 247 [1] Y. Ahn *et al.*, "Review of supercritical CO₂ power cycle technology and current status of research
248 and development," *Nuclear Engineering and Technology*, vol. 47, no. 6, pp. 647–661, Oct. 2015,
249 doi: 10.1016/J.NET.2015.06.009.
- 250 [2] V. Dostal, P. Hejzlar, and M. J. Driscoll, "The Supercritical Carbon Dioxide Power Cycle: Comparison
251 to Other Advanced Power Cycles," *Nucl Technol*, vol. 154, no. 3, pp. 283–301, Jun. 2006, doi:
252 10.13182/NT06-A3734.
- 253 [3] S. M. Besarati and D. Y. Goswami, "Supercritical CO₂ and other advanced power cycles for
254 concentrating solar thermal (CST) systems," *Advances in Concentrating Solar Thermal Research
255 and Technology*, pp. 157–178, Jan. 2017, doi: 10.1016/B978-0-08-100516-3.00008-3.
- 256 [4] E. G. Feher, "The supercritical thermodynamic power cycle," *Energy Conversion*, vol. 8, no. 2, pp.
257 85–90, Sep. 1968, doi: 10.1016/0013-7480(68)90105-8.
- 258 [5] G. Angelino, "Carbon dioxide condensation cycles for power production," *J Eng Gas Turbine
259 Power*, vol. 90, no. 3, pp. 287–295, 1968, doi: 10.1115/1.3609190.
- 260 [6] Y. M. Kim, C. G. Kim, and D. Favrat, "Transcritical or supercritical CO₂ cycles using both low- and
261 high-temperature heat sources," *Energy*, vol. 43, no. 1, pp. 402–415, 2012, doi:
262 <https://doi.org/10.1016/j.energy.2012.03.076>.
- 263 [7] S. A. Wright, R. F. Radel, T. M. Conboy, and G. E. Rochau. (2011). Modeling and experimental
264 results for condensing supercritical CO₂ power cycles, doi: 10.2172/1030354.
- 265 [8] J. Sarkar and S. Bhattacharyya, "Optimization of recompression S-CO₂ power cycle with
266 reheating," *Energy Convers Manag*, vol. 50, no. 8, pp. 1939–1945, Aug. 2009, doi:
267 10.1016/J.ENCONMAN.2009.04.015.
- 268 [9] S. A. Wright, R. F. Radel, M. E. Vernon, P. S. Pickard, and G. E. Rochau. (2010). "Operation and
269 analysis of a supercritical CO₂ Brayton cycle." doi: 10.2172/984129.

270

271 **Tables**

272

Parameters	Reference Model	Thermal Adapter Model
High Pressure	20 MPa	20 MPa
T_{Source}	650°C	650°C
HTR Approach Temp	4°C	4°C
LTR Approach Temp	3.5°C	3.5°C
Total Mass Flow	1 Kg/s	1 Kg/s
Split Ratio (To Heat Rejection)	0.604	0.544
$T_{\text{Rejection}} (T_{\text{Rej}})$	15°C – 45°C	15°C – 45°C
$\eta_{\text{ Main-compressor}}$	0.85	-
$\eta_{\text{ Pump}}$	-	0.85
$\eta_{\text{ Re-compressor}}$	0.87	0.87
$\eta_{\text{ Turbine}}$	0.9	0.9
$\eta_{\text{ Propane Compressor}}$	-	0.9
CO ₂ -Propane HEX Approach Temp	-	3°C
CO ₂ Condensation Temp	-	15°C
Pressure Drop	5% of total	5% of total

Table 1: Model Parameters

273

274

275

276

277

Component	Formula
Exergy flow in	$\Phi_{in} = Q_{in} * (1 - \frac{T_o}{T_{source}})$
Turbomachinery	$i_{Turbo} = \dot{m} * T_o * (s_{out} - s_{in})$
Valve	$i_{valve} = \dot{m} * T_o * (s_{out} - s_{in})$
Heat Exchanger	$i_{HEX} = T_o * [\dot{m}_h(s_{h,out} - s_{h,in}) + \dot{m}_c(s_{c,out} - s_{c,in})]$
Heater	$i_{heater} = T_o * [\dot{m} * (s_{out} - s_{in}) - \frac{Q_{in}}{T_{source}}]$
Cooler	$i_{cooler} = T_o * [\dot{m} * (s_{out} - s_{in}) + \frac{Q_{out}}{T_{sink}}]$

Table 2: Exergy Analysis Calculations

279 **Figures**

280

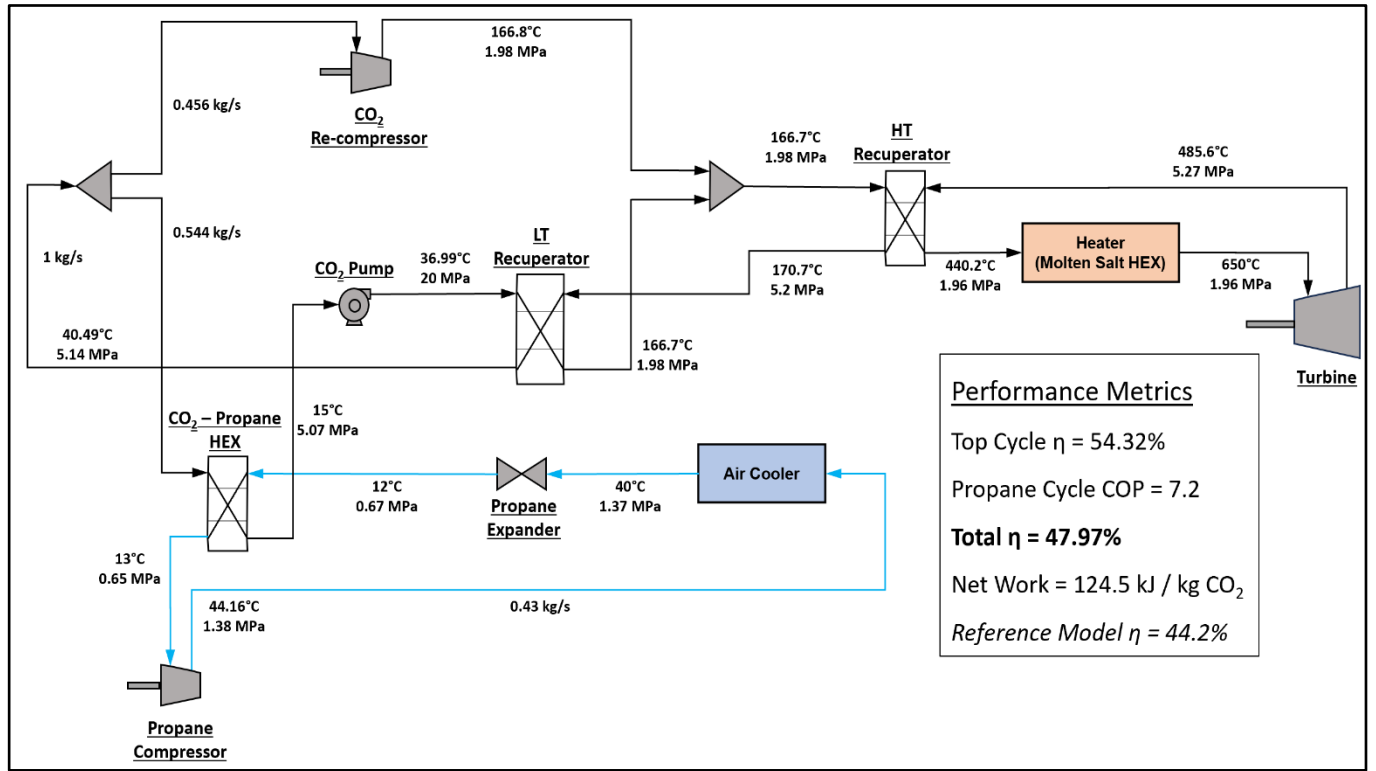


Figure 1: Model Schematic ($T_{Rej} = 40\text{ C}$)

281

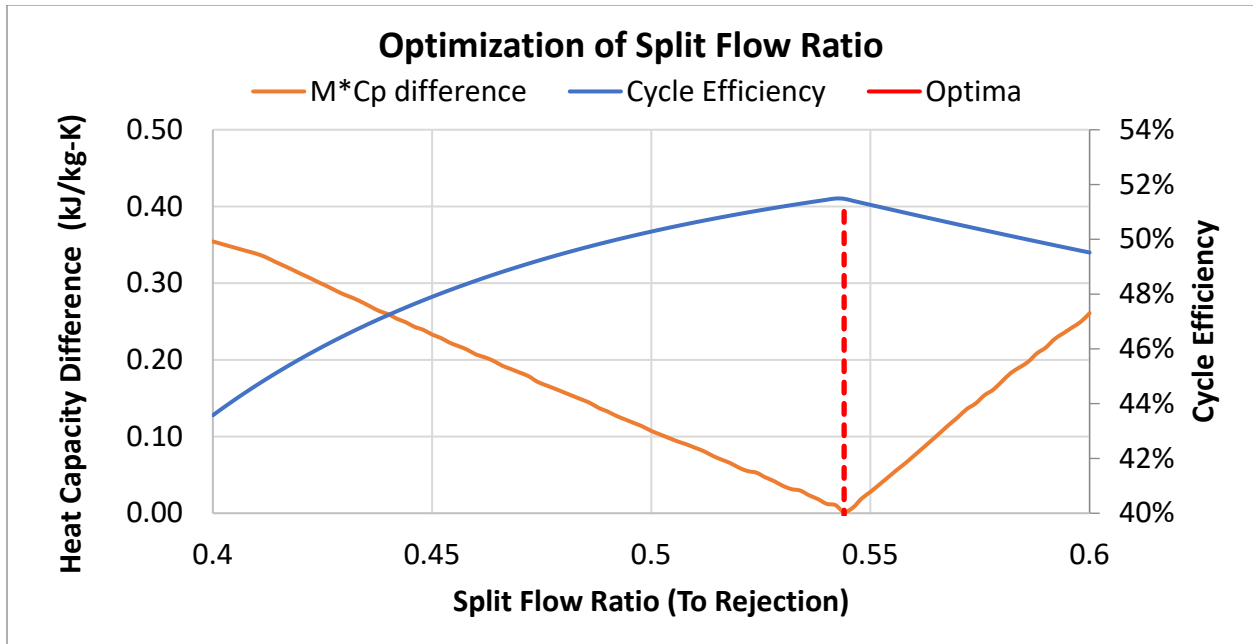
282

283

284

285

286



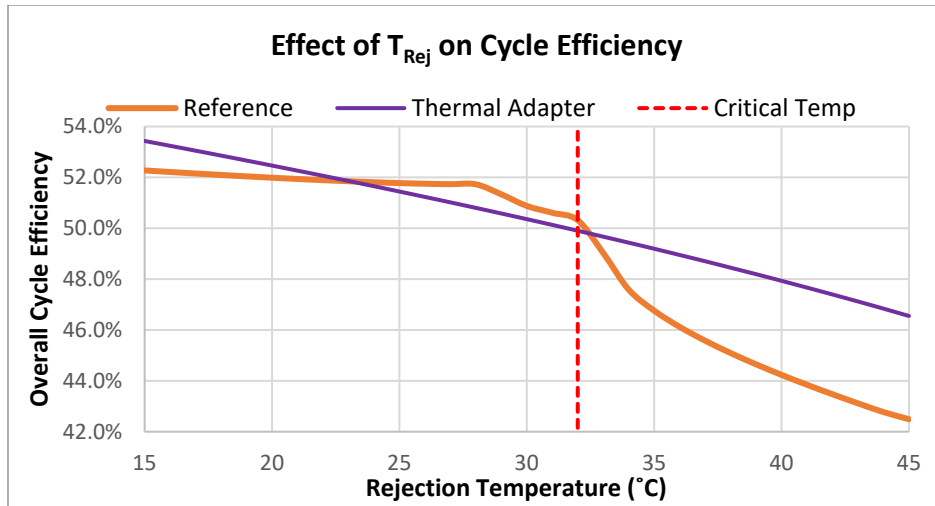
287

288

Figure 2: Split Flow Ratio Optimization

289

290



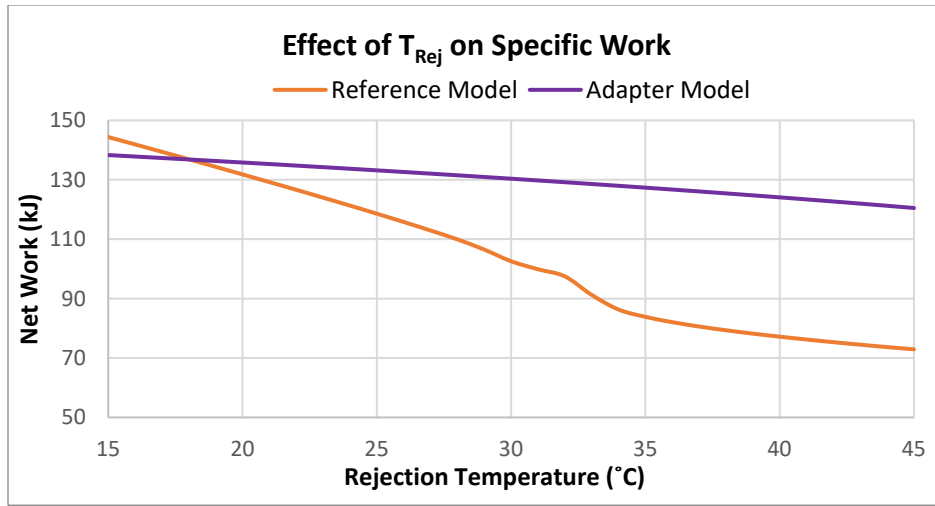
291

292

293

Figure 3: Efficiency v/s T_{Rej}

294



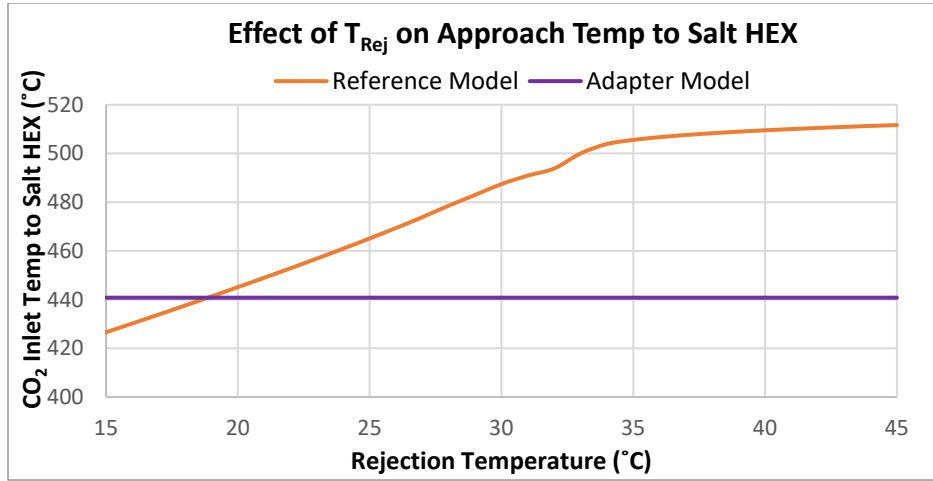
295

296

297

Figure 4: Specific Work v/s T_{Rej}

298



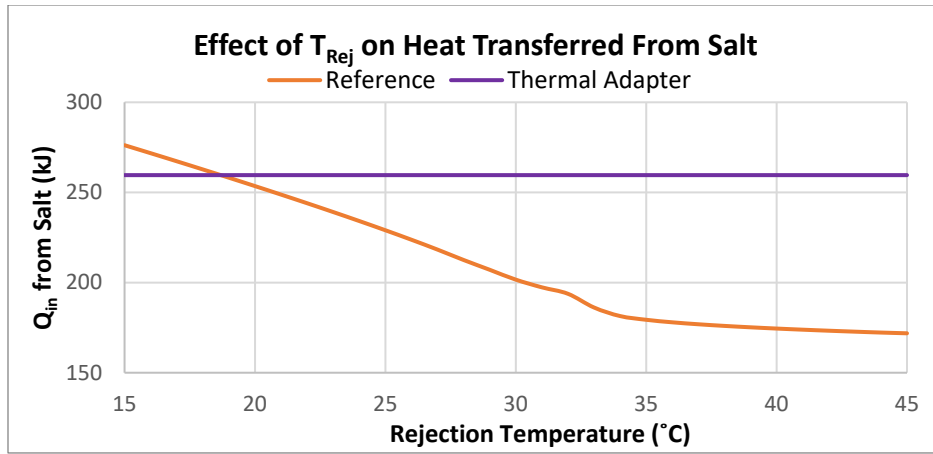
299

300

301

Figure 5: CO₂ Approach temp to Salt HEX v/s T_{Rej}

302



303

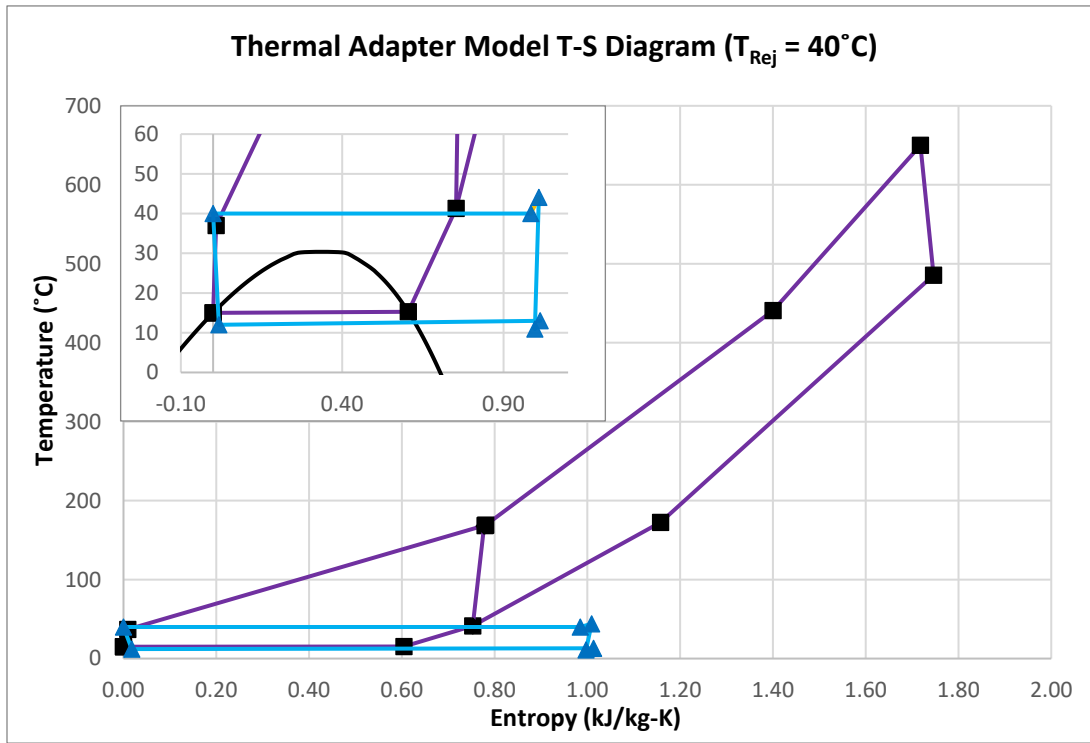
304

Figure 6: Heat duty absorbed from Salt v/s T_{Rej}

305

306

307



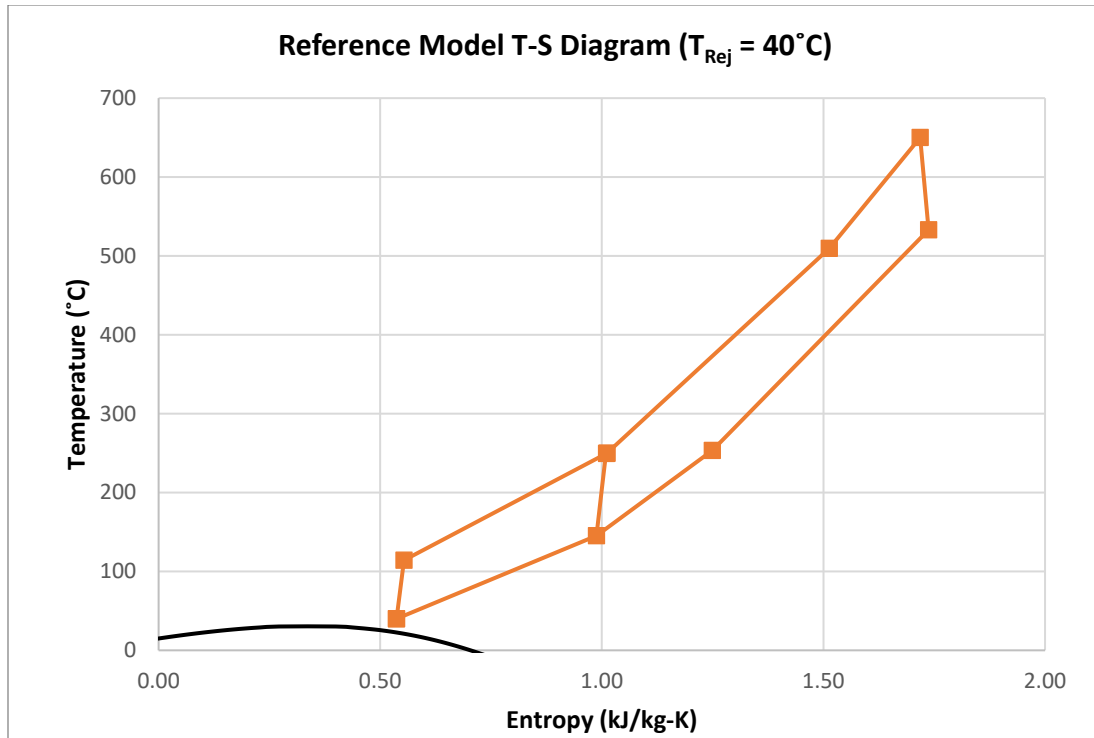
308

309

310

Figure 7: T-S Diagram of Thermal Adapter Model

311

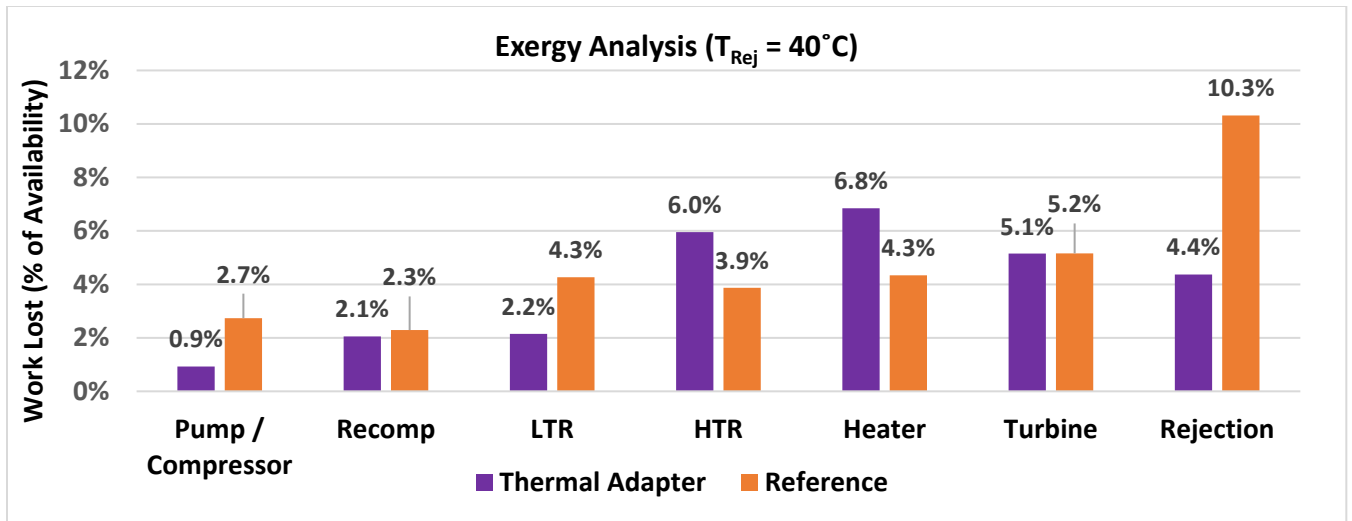


312

313

314

Figure 8: T-S Diagram of Reference Model

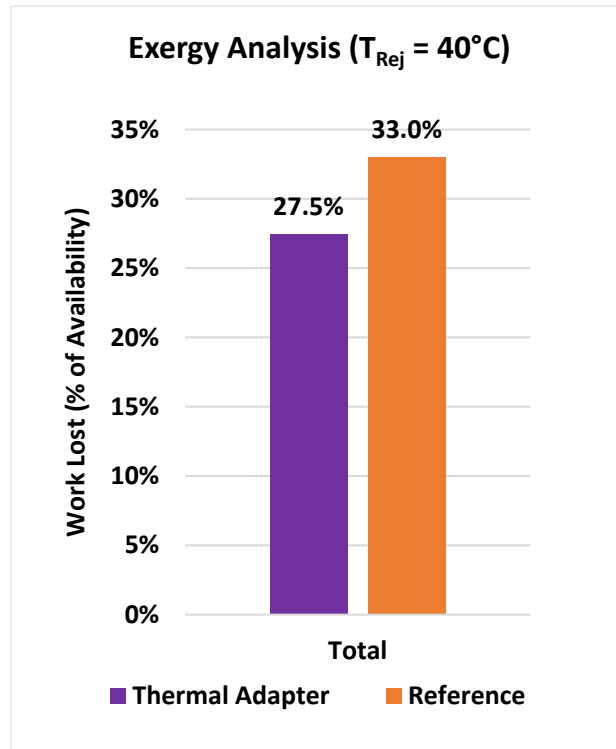


315

316

317

Figure 9: Exergy Analysis Results



318

319

320

321

322

323

324

325

Figure 10: Exergy Analysis Results (Total)

326 **Appendix A**

327 Analysis of Ideal and Real Heat Engine – Refrigerator Combined Cycle

328 Jon D. McWhirter, Ph.D., P.E.

329 TerraPower, LLC

330 **1 PURPOSE**

331 This work is to demonstrate the thermodynamic feasibility of a combined cycle with sub-ambient heat
332 rejection from a power cycle to a refrigeration system then ultimate heat rejection above the ambient
333 temperature.

334 **2 NOMENCLATURE**

335 η - efficiency

336 T - temperature

337 Q - heat

338 β - coefficient of performance of refrigerator

339 W - work

340 **3 SUMMARY**

341 The Supercritical CO₂ Brayton cycle with Recuperation and Recompression, being investigated by
342 various parties for application to nuclear plants but may suffer from issues related to the instability of
343 operations in the vicinity of the critical point. The ambient temperatures worldwide vary enough to
344 be above and below the CO₂ critical point of 31.1 °C. A combined cycle with Supercritical CO₂ power
345 cycle rejecting heat to a refrigeration cycle is a potential scheme to maintain stable operating points
346 of the CO₂ while having the ‘bottoming’ refrigeration cycle as a ‘thermal adapter’ to absorb variations
347 in the ambient temperature while maintaining fixed CO₂ state points. To demonstrate that there is
348 no violation of the Second Law of Thermodynamics, it is shown that the overall cycle efficiency does
349 not exceed the Carnot Cycle efficiency for a heat engine operating between a high temperature
350 reservoir and the ambient temperature. It is emphasized that the combined cycle efficiency is not
351 guaranteed to be higher than the stand-alone Brayton cycle.

352 **4 BACKGROUND**

353 The Supercritical CO₂ Brayton cycle with Recuperation and Recompression is being investigated by
354 various parties for application to nuclear plants (Dostal, Driscoll, & Hejzlar, 2004). However, the
355 potential difficulty of maintaining stability in the vicinity of the critical point, 31.1 °C, and the strong
356 variation of the cycle performance as the ambient temperature changes, are complications, though
357 some analyses show stable operation of a compressor in the vicinity of the critical point (Wright,
358 Radel, Vernon, Rochau, & Pickard, 2010). The notion of using a refrigeration cycle to create a stable
359 *cold space* for the power cycle heat to be rejected to is proposed as a method of avoiding the issues
360 associated with the critical point and the ambient temperature variation by having the heat sink
361 sufficiently cool so that the CO₂ can be readily condensed regardless of the ambient temperature. To
362 demonstrate that there is no violation of the Second Law of Thermodynamics, ideal efficiencies (i.e.,
363 Carnot) for the heat engine and the refrigerator are employed to arrive at equations demonstrating
364 the efficiency for such a scheme relative to the Carnot efficiency.

365 **5 ANALYSIS**

366 5.1 Combined cycle description

367 A schematic of the system is shown below in Figure 1. Four thermal reservoirs are involved: the high
368 temperature Heat Source at T_H , the ultimate Heat Sink at $T_{AMBIENT}$, the heat rejection temperature for
369 the heat engine T_L , and the cold space temperature for the refrigerator, $T_{L'}$. For heat to flow from the
370 heat engine heat sink temperature to the refrigerator cold space, a temperature difference, $DT = T_L -$
371 $T_{L'} \geq 0$, is introduced. Heat from a high temperature reservoir is converted to work in a Carnot cycle
372 with the rejected heat discharged to heat engine heat sink; this heat in its entirety then flows to a
373 colder temperature region generated by the refrigeration cycle. The Carnot refrigerator then takes
374 this quantity of heat, and with a portion of the work produced by the heat engine, compresses the
375 refrigerant to a temperature at ambient; the ultimate heat rejected must include the contribution due
376 to the refrigeration compressor work. The combined cycle net work, W_{NET} is then the heat engine
377 work W_{HE} less the work required for the refrigeration compressor, W_{REF} .

378 5.2 Ideal Analysis

379 Beginning with a standard Carnot Heat Engine (Van Wylen & Sonntag, 1986) operating between a
380 thermal reservoir heat source at T_H and a thermal reservoir heat sink at T_L , we have:

381

$$\eta_{CARNOT} = 1 - \frac{T_L}{T_H} \quad (7-1)$$

382

383 so

$$W_{HE} = Q_H \left(1 - \frac{T_L}{T_H}\right) \quad (7-2)$$

384

385 And, by reversibility and the definition of absolute temperature,

$$\frac{Q_H}{T_H} = \frac{Q_L}{T_L} \quad (7-3)$$

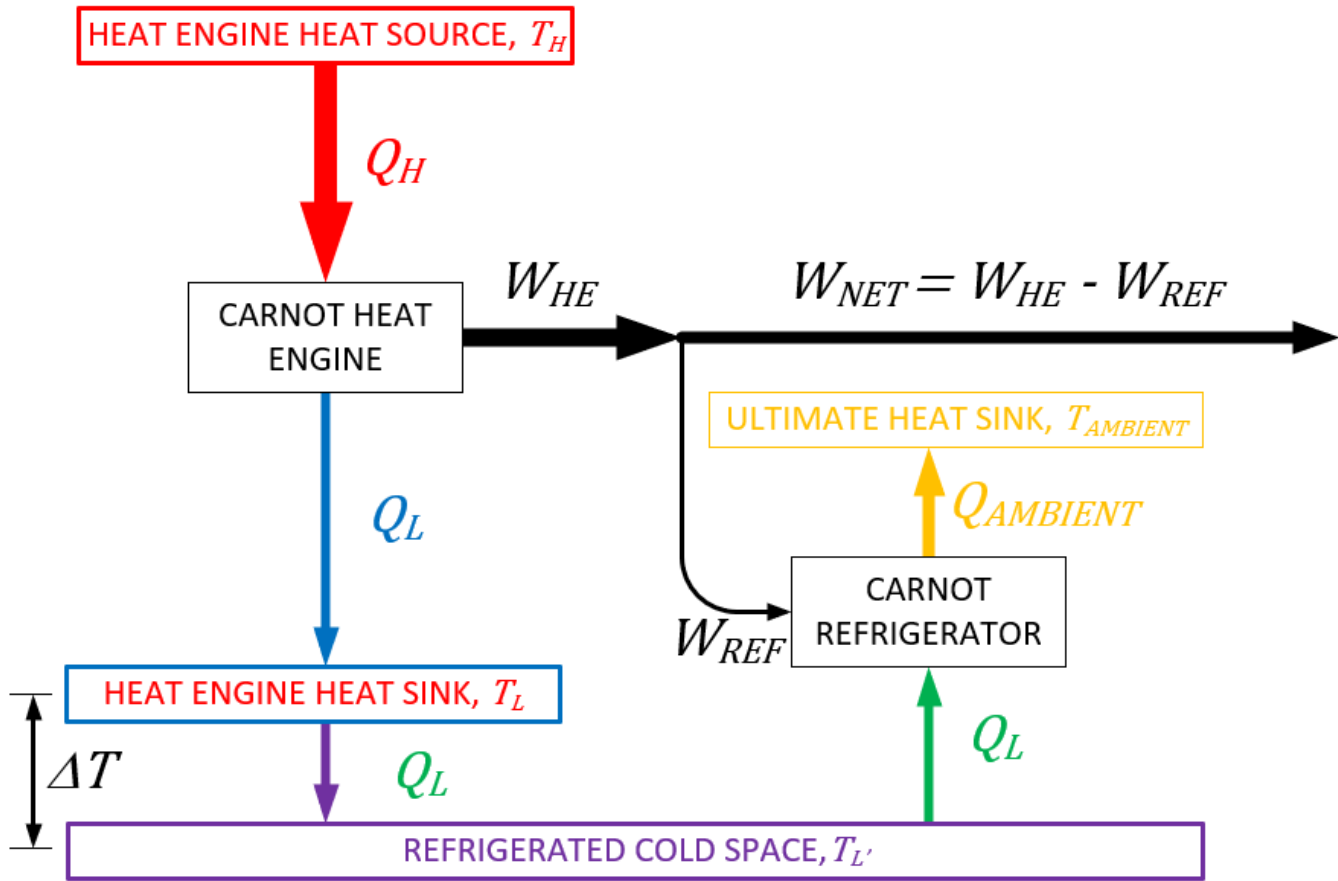
386

387 Thus

$$W_{HE} = Q_L \frac{T_H}{T_L} \left(1 - \frac{T_L}{T_H}\right) = Q_L \left(\frac{T_H}{T_L} - 1\right) \quad (7-4)$$

388

389 Next we introduce a Carnot Refrigerator absorbing heat at $T_{L'}$ and rejecting heat to the ambient at
390 $T_{AMBIENT}$.



391
 392 Figure 11: Schematic of the Combined Cycle

393
 394 The Coefficient of Performance of the Carnot Refrigerator (Van Wylen & Sonntag, 1986) is

$$COP_{CARNOT} = \beta_{CARNOT} = \frac{Q_L}{W_{REF}} = \frac{Q_L}{Q_{AMBIENT} - Q_L} = \frac{T_{L'}}{T_{AMBIENT} - T_{L'}} \quad (7-5)$$

395
 396 So

$$W_{REF} = Q_L \left(\frac{T_{AMBIENT} - T_{L'}}{T_{L'}} \right) = Q_L \left(\frac{T_{AMBIENT}}{T_{L'}} - 1 \right) \quad (7-6)$$

397
 398 Hence

$$W_{NET} = W_{HE} - W_{REF} = Q_L \left(\frac{T_H}{T_L} - 1 \right) - Q_L \left(\frac{T_{AMBIENT}}{T_{L'}} - 1 \right) \quad (7-7)$$

399

400

401 The goal is to find the overall efficiency, $\eta_{OVERALL}$, as a function of the thermal reservoir and ambient
 402 temperatures and the difference between the two low-temperature reservoirs, DT . Since

$$W_{NET} = W_{HE} - W_{REF} = Q_L \left\{ \left(\frac{T_H}{T_L} - 1 \right) - \left(\frac{T_{AMBIENT}}{T_{L'}} - 1 \right) \right\} = Q_L \left\{ \frac{T_H}{T_L} - \frac{T_{AMBIENT}}{T_{L'}} \right\} \quad (7-8)$$

$$W_{NET} = Q_H \frac{T_L}{T_H} \left\{ \frac{T_H}{T_L} - \frac{T_{AMBIENT}}{T_{L'}} \right\} = Q_H \left\{ 1 - \frac{T_L}{T_H} \frac{T_{AMBIENT}}{T_{L'}} \right\}$$

$$\therefore \eta_{OVERALL} = \frac{W_{NET}}{Q_H} = \left\{ 1 - \frac{T_L}{T_H} \frac{T_{AMBIENT}}{T_{L'}} \right\}$$

403

404 With the power cycle heat rejection temperature slightly higher than the refrigeration cycle heat
 405 removal temperature,

406

$$T_L = T_{L'} + \Delta T, \Delta T > 0 \quad (7-9)$$

407

408 there results:

$$\eta_{OVERALL} = \frac{W_{NET}}{Q_H} = 1 - \left(\frac{T_{L'} + \Delta T}{T_H} \right) \left(\frac{T_{AMBIENT}}{T_{L'}} \right)$$

$$\therefore \eta_{OVERALL} = 1 - \left(\frac{T_{L'} + \Delta T}{T_{L'}} \right) \left(\frac{T_{AMBIENT}}{T_H} \right) \quad (7-10)$$

409

410 Now for some limiting cases. Note that for $DT = 0$, we arrive at the expected overall efficiency for an
 411 ideal heat engine operating between the high temperature reservoir and the ambient, η_{CARNOT} :

$$\Delta T = 0 \rightarrow \eta_{OVERALL} = 1 - \left(\frac{T_{L'}}{T_{L'}} \right) \left(\frac{T_{AMBIENT}}{T_H} \right) = 1 - \frac{T_{AMBIENT}}{T_H} = \eta_{CARNOT} \quad (7-11)$$

412

413 for a cycle between T_H and $T_{AMBIENT}$. Now, as $DT > 0$ (by definition), the overall efficiency is less than
 414 the Carnot efficiency – the numerator of the subtrahend in Eq. (7-11) is larger than that of the ideal
 415 case, so the difference is smaller. Hence, the Second Law of Thermodynamics remains observed since
 416 the Carnot cycle efficiency is not exceeded. Some other observations of the mathematics can be
 417 demonstrated more readily by looking at the ratio of the Carnot efficiency and the Overall Efficiency:

$$\frac{\eta_{OVERALL,\Delta T \neq 0}}{\eta_{CARNOT}} = \frac{1 - \left(\frac{T_{L'} + \Delta T}{T_{L'}}\right) \left(\frac{T_{AMBIENT}}{T_H}\right)}{1 - \frac{T_{AMBIENT}}{T_H}} = \frac{1 - \left(\frac{T_{AMBIENT}}{T_H}\right) - \left(\frac{\Delta T}{T_{L'}}\right) \left(\frac{T_{AMBIENT}}{T_H}\right)}{1 - \frac{T_{AMBIENT}}{T_H}}$$

$$\frac{\eta_{OVERALL,\Delta T \neq 0}}{\eta_{CARNOT}} = \frac{(T_H - T_{AMBIENT}) - T_{AMBIENT} \left(\frac{\Delta T}{T_{L'}}\right)}{(T_H - T_{AMBIENT})}$$

$$\therefore \frac{\eta_{OVERALL,\Delta T \neq 0}}{\eta_{CARNOT}} = 1 - \frac{T_{AMBIENT} \left(\frac{\Delta T}{T_{L'}}\right)}{T_H - T_{AMBIENT}} < 1, \forall \Delta T > 0 \quad (7-12)$$

418 So as derived above, the Carnot efficiency is approached as $\Delta T \rightarrow 0$. Also, the relative deviation from
419 Carnot efficiency rises as $T_{AMBIENT}$ rises, a reasonable expectation.

420 6 REAL CYCLES

421 The objective of the real cycle analysis is to express the efficiency of an overall cycle in terms of the
422 performance metrics of the heat engine and the refrigeration system. For a real heat engine with
423 efficiency h and a refrigeration cycle with coefficient of performance b , with appropriate
424 temperatures to foster the heat flows from higher to lower temperature, the work produced by the
425 heat engine, for a heat input Q_H , is $W_{HE} = h Q_H$. The rejected heat is, therefore,

$$426 \quad \quad \quad Q_L = (1 - \eta)Q_H \quad (7-13)$$

427
428 which is the heat that must be *lifted and rejected* by the refrigeration cycle. The coefficient of
429 performance, β , for the refrigeration cycle is

$$430 \quad \quad \quad \beta = Q_L/W_{REF} \Rightarrow Q_L = \beta W_{REF} \quad (7-14)$$

431 substituting for Q_L yields

$$432 \quad \quad \quad W_{REF} = (1 - \eta)Q_H/\beta \quad (7-15)$$

433

434

435 The net work for the system is the heat engine work less the refrigeration cycle work, or

436

$$W_{NET} = W_{HE} - W_{REF} = \eta Q_H - \frac{(1 - \eta)Q_H}{\beta} = Q_H(\eta - (1 - \eta)/\beta) \quad (7-16)$$

437 The overall efficiency is found by dividing the net work by the heat input, or

438

$$\eta_{OVERALL} = \eta - (1 - \eta)/\beta \quad (7-17)$$

439

440 **7 CONCLUSIONS**

441 Ideal and real analyses have demonstrated that such a combined cycle with a *cold space* operating
442 below ambient temperature, created by a refrigeration cycle rejecting heat to ambient, follows
443 expected behavior mathematically. However, the test of whether this results in a gain over the
444 'original' Brayton cycle rejecting heat directly to ambient alone is not evaluated herein. That analysis
445 is more complicated mathematically and must be realistically evaluated numerically with real
446 thermodynamic state point data.

447 **8 BIBLIOGRAPHY**

- 448 Dostal, V., Driscoll, M., & Hejzlar, P. (2004). *A Supercritical Carbon Dioxide Cycle for Next Generation*
449 *Nuclear Reactors*. Cambridge, Massachusetts : MIT.
- 450 Van Wylen, G., & Sonntag, E. (1986). *Fundamentals of Classical Thermodynamics, 3rd edition*. John Wiley
451 & Sons.
- 452 Wright, S., Radel, R., Vernon, M., Rochau, G., & Pickard, P. (2010). *Operation and Analysis of a Supercritical*
453 *CO2 Brayton Cycle, SAND2010-0171*. Sandia National Laboratory.

454

455

456